

THE FUNDAMENTAL SOLUTIONS TO THE CENTRAL SYMMETRIC TIME-FRACTIONAL HEAT CONDUCTION EQUATION WITH HEAT ABSORPTION

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Received: 8 March 2017; accepted: 12 June 2017

Abstract. The time-fractional heat conduction equation with heat absorption proportional to temperature is considered in the case of central symmetry. The fundamental solutions to the Cauchy problem and to the source problem are obtained using the integral transform technique. The numerical results are presented graphically.

MSC 2010: 35K05, 35R11, 26A33, 44A10, 42A38

Keywords: non-Fourier heat conduction, Caputo fractional derivative, heat absorption, Laplace integral transform, Fourier transform, Mittag-Leffler function

1. Introduction

The classical heat conduction is based on the standard Fourier law and the parabolic heat conduction equation. The time-nonlocal dependence of the heat flux and the temperature gradient with the “long-tail” power kernel [1-4] can be interpreted in terms of fractional integrals and derivatives and results in the time-fractional heat conduction equation

$$\frac{\partial^\alpha T}{\partial t^\alpha} = a \Delta T, \quad 0 < \alpha \leq 2, \quad (1)$$

where T is a temperature, t denotes time, Δ stands for the Laplace operator, a is an analogue of the thermal diffusivity coefficient. The Caputo fractional derivative is defined as [5-7]

$$\frac{d^\alpha T}{dt^\alpha} = \frac{1}{\Gamma(n-\alpha)} \int_0^t (t-\tau)^{n-\alpha-1} \frac{d^n T(\tau)}{d\tau^n} d\tau, \quad n-1 < \alpha < n, \quad (2)$$

with $\Gamma(x)$ being the gamma function.

Fractional calculus (the theory of integrals and derivatives of non-integer order) provides the appropriate mathematical tool for description of many phenomena in physics, chemistry, biology, and engineering [8-15].

If volume heat absorption proportional to temperature occurs in a body, then instead of (1) we get

$$\frac{\partial^\alpha T}{\partial t^\alpha} = a \Delta T - bT, \quad 0 < \alpha \leq 2, \quad (3)$$

where the values of the coefficient $b > 0$ and $b < 0$ correspond to heat absorption and heat release, respectively. The classical heat conduction equation with the additional term proportional to temperature was considered in [16-18]. Similar equations appear in the theory of bio-heat transfer [19] and in the survival probability [20]. Mathematical and physical aspects of fractional heat conduction equation with heat absorption were studied in the literature in the case of one Cartesian coordinate in [21-24]. In the present paper, we study the fundamental solutions to the Cauchy problem and to the source problem for equation (3) in spherical coordinates in the case of central symmetry. The obtained solutions generalize the results of the paper [25], where the case $b = 0$ was considered.

2. The fundamental solution to the Cauchy problem

We consider the time-fractional heat conduction equation with one spatial variable in spherical coordinate system

$$\frac{\partial^\alpha T(r,t)}{\partial t^\alpha} = a \left(\frac{\partial^2 T(r,t)}{\partial r^2} + \frac{2}{r} \frac{\partial T(r,t)}{\partial r} \right) - bT(r,t), \quad (4)$$

where $0 < r < \infty$, $0 < t < \infty$, $0 < \alpha \leq 2$, $a > 0$.

Equation (4) is considered under initial conditions

$$T(r,0) = \frac{P_0}{4\pi r^2} \delta(r-R), \quad 0 < \alpha \leq 2, \quad (5)$$

$$\frac{\partial T(r,0)}{\partial t} = 0, \quad 1 < \alpha \leq 2, \quad (6)$$

with $\delta(r)$ being the Dirac delta function. For the sake of convenience and to obtain the nondimensional quantities used in calculations, we have introduced the constant multiplier p_0 in equation (5).

The zero condition at infinity is also imposed:

$$\lim_{r \rightarrow \infty} T(r, t) = 0. \tag{7}$$

To solve the Cauchy problem under consideration we use the integral transform technique. The Laplace transform with respect to the time t is defined as

$$L\{T(t)\} = T^*(s) = \int_0^\infty e^{-st} T(t) dt \tag{8}$$

with the inverse carried out according to the Fourier-Mellin formula

$$L^{-1}\{T^*(s)\} = T(t) = \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} e^{st} T^*(s) ds, \tag{9}$$

where s denotes the transform variable, c is a positive fixed number such that all the singularities of $T^*(s)$ lie to the left of the vertical line $\text{Re } s = c$.

For the Laplace transform rule, the Caputo fractional derivative requires the knowledge of the initial values of the function and its integer derivatives of the order $k = 1, 2, \dots, n - 1$:

$$L\left\{\frac{d^\alpha T(t)}{dt^\alpha}\right\} = s^\alpha T^*(s) - \sum_{k=0}^{n-1} s^{\alpha-1-k} T^{(k)}(0^+), \quad n-1 < \alpha \leq n. \tag{10}$$

Applying the Laplace transform to equation (4) and taking into account the rule (10) with the initial conditions (5) and (6) gives

$$s^\alpha T^*(r, s) - s^{\alpha-1} \frac{P_0}{4\pi r^2} \delta(r - R) = a \left(\frac{\partial^2 T^*(r, s)}{\partial r^2} + \frac{2}{r} \frac{\partial T^*(r, s)}{\partial r} \right) - b T^*(r, s). \tag{11}$$

Next, we use the Fourier transform with respect to the spatial coordinate r in the case of spherical symmetry [4, 26]:

$$F\{T(r)\} = \tilde{T}(\rho) = \sqrt{\frac{2}{\pi}} \int_0^\infty r T(r) \frac{\sin(r\rho)}{\rho} dr, \tag{12}$$

$$F^{-1}\{\tilde{T}(\rho)\} = T(r) = \sqrt{\frac{2}{\pi}} \int_0^\infty \rho \tilde{T}(\rho) \frac{\sin(r\rho)}{r} d\rho. \tag{13}$$

$$F\left\{\frac{d^2 T(r)}{dr^2} + \frac{2}{r} \frac{dT(r)}{dr}\right\} = -\rho^2 \tilde{T}(\rho). \quad (14)$$

Usually the Fourier transform (12)-(13) is used under the assumption of boundedness of $T(r)$ at the origin (see, e.g., [18]); sometimes this assumption is substituted by less restrictive condition prescribing a type of singularity of the function at $r = 0$ (see, for instance, [27]).

Application of the Fourier transform (12) and formula (14) to equation (11) leads to

$$\tilde{T}^*(\rho, s) = \frac{p_0}{4\pi} \sqrt{\frac{2}{\pi}} \frac{\sin(\rho R)}{\rho R} \frac{s^{\alpha-1}}{s^\alpha + a\rho^2 + b}. \quad (15)$$

Inversion of the integral transforms results in the solution:

$$T(r, t) = \frac{p_0}{2\pi^2 R r} \int_0^\infty \sin(\rho R) \sin(\rho r) E_\alpha[-(a\rho^2 + b)t^\alpha] d\rho, \quad (16)$$

where $E_\alpha(z)$ is the Mittag-Leffler function in one parameter α [5-7]

$$E_\alpha(z) = \sum_{n=0}^{\infty} \frac{z^n}{\Gamma(\alpha n + 1)}, \quad \alpha > 0, \quad z \in C, \quad (17)$$

and the following formula has been used:

$$L^{-1}\left\{\frac{s^{\alpha-1}}{s^\alpha + g}\right\} = E_\alpha(-gt^\alpha). \quad (18)$$

Using the nondimensional quantities

$$\bar{r} = \frac{r}{R}, \quad \bar{\rho} = R\rho, \quad \kappa = \frac{\sqrt{a} t^{\alpha/2}}{R}, \quad \bar{b} = bt^\alpha, \quad \bar{T} = \frac{4\pi^2 R^3}{p_0} T, \quad (19)$$

one obtains the following solution:

$$\bar{T}(\bar{r}, \kappa) = \frac{2}{\bar{r}} \int_0^\infty \sin(\bar{\rho}) \sin(\bar{r}\bar{\rho}) E_\alpha[-(\kappa^2 \bar{\rho}^2 + \bar{b})] d\bar{\rho}. \quad (20)$$

For $\alpha = 1$ the Mittag-Leffler function $E_1(z) = e^z$, and from (20) we get

$$\bar{T}(\bar{r}, \kappa) = \frac{1}{\bar{r}} e^{-\bar{b}} \left\{ \int_0^\infty e^{-\kappa^2 \bar{\rho}^2} \cos[\bar{\rho}(\bar{r} - 1)] d\bar{\rho} - \int_0^\infty e^{-\kappa^2 \bar{\rho}^2} \cos[\bar{\rho}(\bar{r} + 1)] d\bar{\rho} \right\}. \quad (21)$$

Taking into account that [28]

$$\int_0^\infty \exp(-p^2 x^2) \cos(qx) dx = \frac{\sqrt{\pi}}{2p} \exp\left(-\frac{q^2}{4p^2}\right), \quad p > 0, \quad (22)$$

we arrive at the solution

$$\bar{T}(\bar{r}, \kappa) = \frac{\sqrt{\pi}}{2\kappa\bar{r}} e^{-\bar{b}} \left\{ \exp\left[-\frac{(\bar{r}-1)^2}{4\kappa^2}\right] - \exp\left[-\frac{(\bar{r}+1)^2}{4\kappa^2}\right] \right\}. \quad (23)$$

In the particular case $\alpha = 1/2$, the Mittag-Leffler function $E_{1/2}(-x)$ can be represented as [4]

$$E_{1/2}(-x) = \frac{2}{\sqrt{\pi}} \int_0^\infty \exp(-u^2 - 2ux) du, \quad (24)$$

and the solution has the following form

$$\bar{T}(\bar{r}, \kappa) = \frac{1}{\sqrt{2\bar{r}\kappa}} \int_0^\infty \frac{1}{\sqrt{u}} e^{-u^2 - 2u\bar{b}} \left\{ \exp\left[-\frac{(\bar{r}-1)^2}{8u\kappa^2}\right] - \exp\left[-\frac{(\bar{r}+1)^2}{8u\kappa^2}\right] \right\} du. \quad (25)$$

The results of numerical calculations for $\kappa = 0.25$, different values of \bar{b} and the order of the time-fractional derivative α are shown in Figures 1-5.

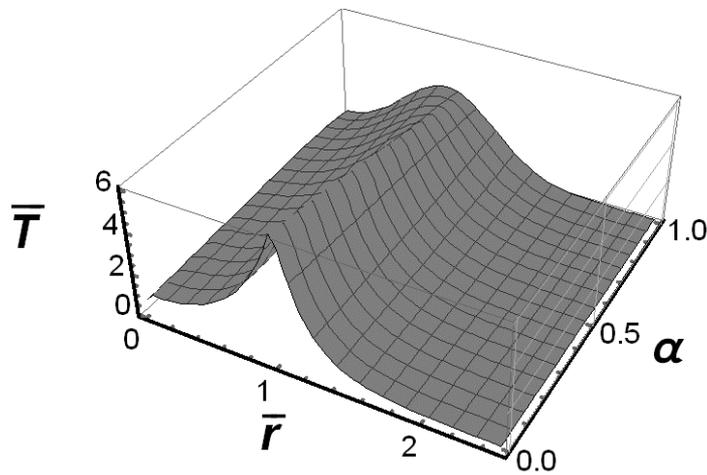


Fig. 1. The fundamental solution to the Cauchy problem for $\bar{b} = 0$ and $0 \leq \alpha \leq 1$

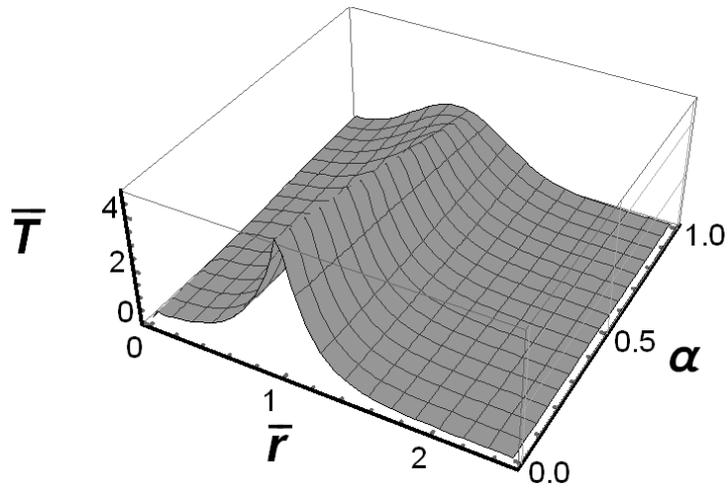


Fig. 2. The fundamental solution to the Cauchy problem for $\bar{b} = 0.5$ and $0 \leq \alpha \leq 1$

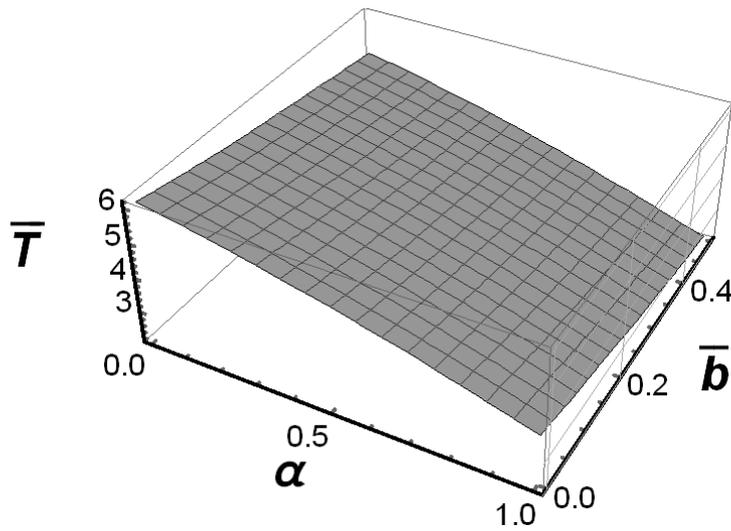


Fig. 3. The fundamental solution to the Cauchy problem for $\bar{r} = 1$, $0 \leq \alpha \leq 1$ and $0 \leq \bar{b} \leq 0.5$

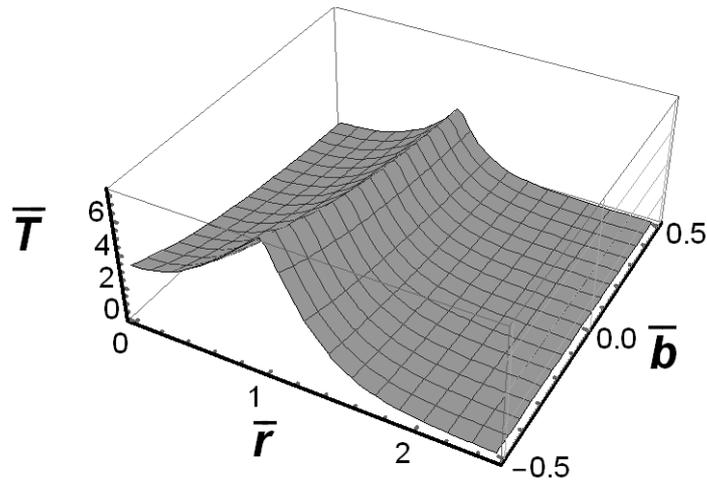


Fig. 4. The fundamental solution to the Cauchy problem for $\alpha = 0.5$ and $-0.5 \leq \bar{b} \leq 0.5$

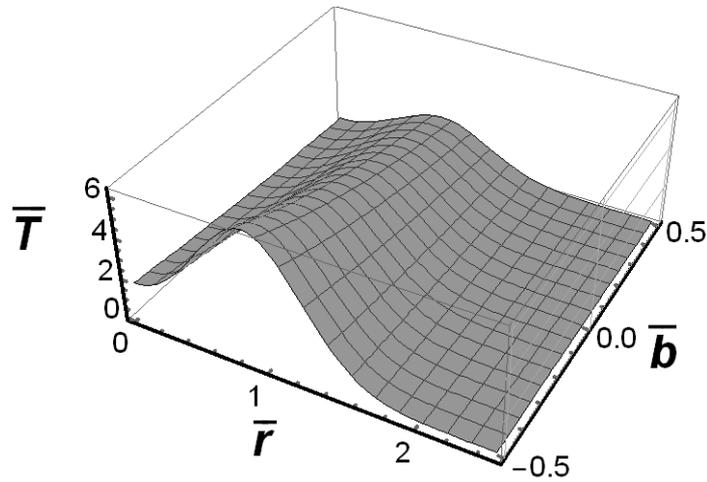


Fig. 5. The fundamental solution to the Cauchy problem for $\alpha = 1$ and $-0.5 \leq \bar{b} \leq 0.5$

3. The fundamental solution to the source problem

Consider the time-fractional heat conduction equation with the source term

$$\frac{\partial^\alpha T(r,t)}{\partial t^\alpha} = a \left(\frac{\partial^2 T(r,t)}{\partial r^2} + \frac{2}{r} \frac{\partial T(r,t)}{\partial r} \right) - bT(r,t) + \frac{q_0}{4\pi r^2} \delta(r-R)\delta(t) \quad (26)$$

under zero initial conditions

$$T(r, 0) = 0, \quad 0 < \alpha \leq 2, \quad (27)$$

$$\frac{\partial T(r, 0)}{\partial t} = 0, \quad 1 < \alpha \leq 2. \quad (28)$$

Using the integral transforms technique, we obtain

$$\tilde{T}^*(\rho, s) = \frac{q_0}{4\pi} \sqrt{\frac{2}{\pi}} \frac{\sin(\rho R)}{\rho R} \frac{1}{s^\alpha + a\rho^2 + b}. \quad (29)$$

As

$$L^{-1}\left\{\frac{1}{s^\alpha + g}\right\} = t^{\alpha-1} E_{\alpha, \alpha}(-gt^\alpha), \quad (30)$$

where $E_{\alpha, \beta}(z)$ is the Mittag-Leffler function in two parameters α and β [5-7]

$$E_{\alpha, \beta}(z) = \sum_{n=0}^{\infty} \frac{z^n}{\Gamma(\alpha n + \beta)}, \quad \alpha > 0, \beta > 0, z \in C, \quad (31)$$

the inverse transforms applied to equation (29) lead to

$$T(r, t) = \frac{q_0}{2\pi^2 R r} t^{\alpha-1} \int_0^\infty \sin(\rho R) \sin(\rho r) E_{\alpha, \alpha}[-(a\rho^2 + b)t^\alpha] d\rho \quad (32)$$

and

$$\bar{T}(\bar{r}, \kappa) = \frac{2}{\bar{r}} \int_0^\infty \sin(\bar{\rho}) \sin(\bar{r}\bar{\rho}) E_{\alpha, \alpha}[-(\kappa^2 \bar{\rho}^2 + \bar{b})] d\bar{\rho}. \quad (33)$$

In this case, the nondimensional temperature is introduced as

$$\bar{T} = \frac{4\pi^2 R^3}{q_0 t^{\alpha-1}} T, \quad (34)$$

and other nondimensional quantities are the same as in (19).

In the case $\alpha = 1/2$ [4]

$$E_{1/2, 1/2}(-x) = \frac{2}{\sqrt{\pi}} \int_0^\infty u \exp(-u^2 - 2ux) du. \quad (35)$$

Taking into account (35), (38) and (25), we arrive at

$$\bar{T}(\bar{r}, \kappa) = \frac{1}{\sqrt{2} \bar{r} \kappa} \int_0^\infty \sqrt{u} e^{-u^2 - 2u\bar{b}} \left[\exp\left(-\frac{(\bar{r}-1)^2}{8u\kappa^2}\right) - \exp\left(-\frac{(\bar{r}+1)^2}{8u\kappa^2}\right) \right] du. \quad (36)$$

Figures 6-8 show the results of numerical calculations according to equations (33) and (36) for $\kappa = 0.25$.

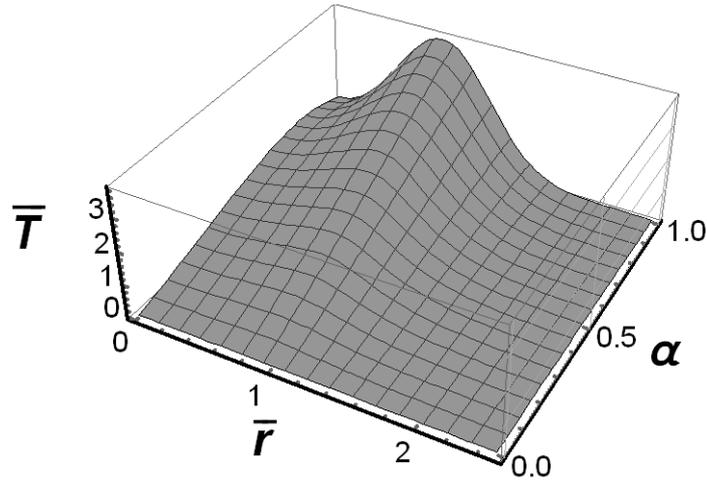


Fig. 6. The fundamental solution to the source problem for $\bar{b} = 0$ and $0 \leq \alpha \leq 1$

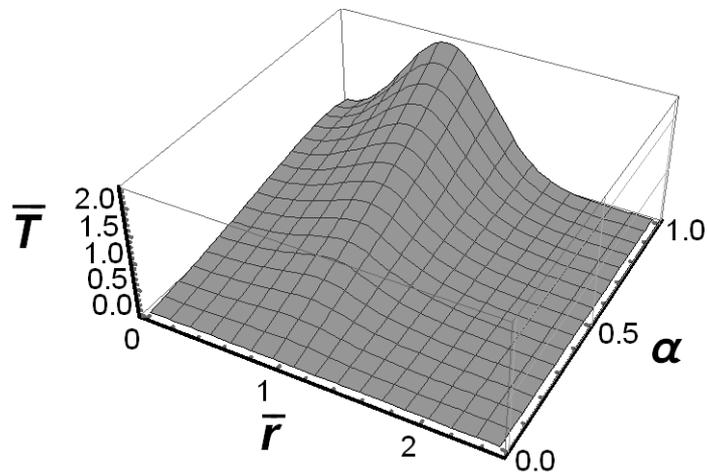


Fig. 7. The fundamental solution to the source problem for $\bar{b} = 0.5$ and $0 \leq \alpha \leq 1$

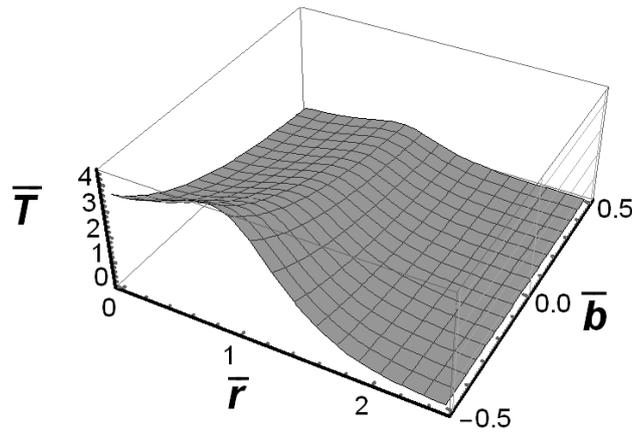


Fig. 8. The fundamental solution to the source problem for $\alpha = 0.5$ and $-0.5 \leq \bar{b} \leq 0.5$

4. Conclusions

We have solved the time-fractional heat conduction equation with the Caputo fractional derivative in the case of one spatial variable in spherical coordinates. The heat absorption is assumed to be proportional to temperature. The fundamental solutions to the Cauchy problem and to the source problem have been studied. It should be noted that in the case of the classical parabolic heat conduction equation ($\alpha = 1$), the fundamental solutions to the Cauchy problem and to the source problem coincide, whereas for $\alpha \neq 1$ they are different. The results of numerical calculations are displayed in figures for different values of the parameter b describing heat absorption and the order of the Caputo fractional derivative. The particular cases of the solutions corresponding to the value $b = 0$ coincide with those obtained in [4, 25]. The influence of the sign change of the parameter b on temperature is easily observable from the figures. To calculate the Mittag-Leffler functions $E_\alpha(-x)$ in (20) and $E_{\alpha,\alpha}(-x)$ in (33), we have used the algorithms suggested in [29]. It should be emphasized that fractional heat conduction and fractional diffusion have the same origin. At the level of individual particle motions the classical diffusion corresponds to Brownian motion with a mean-squared displacement increasing linearly with time. Anomalous diffusion, which is exemplified by a mean-squared displacement with the power-law time dependence and was observed in different media [13, 30-32], is described by equations with fractional derivatives.

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